

Thermodynamic Characterization of Gas Mixtures for Non-Thermal Plasma CO₂ Conversion Applications with Soft-SAFT

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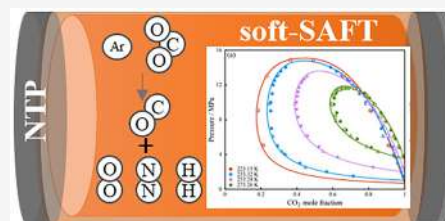
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ABSTRACT: Carbon dioxide (CO₂) transformation into added-value products through non-thermal plasma (NTP) represents a novel technology of interest. The process involves, apart from CO₂, mixtures of different gases such as carbon monoxide (CO), oxygen (O₂), nitrogen (N₂), argon (Ar), and hydrogen (H₂) for subsequent CO₂ methanation. In this work, a preliminary study of the thermodynamic representation of the mixtures relevant in the context of carbon capture, utilization, and storage (CCUS) processes, but focused on the NTP conversion, is presented. The thermodynamic characterization is achieved through the application of the polar soft-statistical associating fluid theory (SAFT) equation of state (EoS), which allows molecular parameterization of pure compounds and the description of mixtures at different conditions of temperature and pressure. An accurate parametrization of all gases is carried out by explicitly considering the quadrupolar nature of CO₂, CO, and N₂. The characterization is then used to describe several single-phase densities, derivative properties, second virial coefficients, and the vapor–liquid equilibrium (VLE) of CO₂ binary mixtures with Ar, O₂, CO, N₂, and H₂, as well as combinations between some of these gases. A parametric analysis of the impact of the binary parameters on the equilibria description is carried out to assess the temperature dependency. The results have overall shown good agreement to experimental data in most conditions using one or two binary parameters. Finally, ternary systems involving CO₂, O₂, Ar, and N₂ have been predicted in good agreement with the experimental data, demonstrating the capacity of the model to evaluate the behavior of multicomponent gas mixtures.



1. INTRODUCTION

The increasing amount of anthropogenic CO₂ emissions has been a major contributor to global warming. According to the very recent 6th Assessment IPCC Report, evidence suggests that around 64% of GHG emissions come from CO₂ originated from fossil fuels and industrial processes.¹ As a consequence, limiting these CO₂ emissions becomes imperative. In this regard, carbon capture, utilization, and storage (CCUS) methods play an important role in net-zero energy systems as they contribute to both reducing and utilizing CO₂ emissions.² In fact, CCUS involves a diverse variety of techniques, developments, and potential applications of high interest. The reader is referred to the recent detailed review of Zhu for further details.³ In particular, CO₂ usage technologies introduce opportunities to partially substitute fossil-fuel raw materials to increase the use of renewable energy as alternatives, as well as to generate income by producing marketable fuel products.

CO₂ usage can be understood as either a direct utilization or a conversion to other chemicals and energy products. On one side, in direct utilization, CO₂ is used as a commercial product in the food and beverage industry or as a solvent in supercritical extraction of natural products, among other industrial applications.⁴ On the other side, CO₂ conversion involves the use of CO₂ as a reactant to be converted into other chemicals or energy products, either through carboxylation reactions in which CO₂ is used to obtain organic

compounds (including carbonates, acrylates, and polymers) or through reduction reactions where the C=O bonds are broken to produce other chemicals (including methane, methanol, urea, syngas, and formic acid).^{5,6} The limitation of CO₂ conversion is the energy intensive nature of these processes due to the high thermodynamic stability of the CO₂ molecule.⁷

CO₂ conversion methods are frequently grouped into: hydrogenation,⁸ electro-catalytic/electrochemical reduction,⁵ enzymatic/biochemical processes,^{5,9,10} photo-reduction,^{5,9,10} or plasma processes, such as non-thermal plasma (NTP).^{3,11–14} The resulting products generated in each of these processes are varied but primarily include: hydrocarbons, hydrogen, and oxygenates via catalytic conversion; and carbon monoxide, hydrocarbons, syngas, and oxygenates from plasma processes.^{7,14} These products have the potential to replace petrochemical feedstocks as the process results in added-value chemicals and fuels.¹⁵

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Among these many CO₂ utilization technologies, a conversion method of growing interest is NTP reduction.^{11–14,16} NTP is generated by applying a potential difference between two electrodes, which are inserted in a reactor filled with gas.¹⁶ When electrons are accelerated by the electric field generated, they collide toward gas molecules, resulting in ionization, excitation, and dissociation. The dissociation collisions create radicals with enough energy input to break the OC=O double bond at atmospheric conditions. It is under this principle that gas conversion applications occur. This technology can be operated at ambient temperature while still generating highly active species and electrons, representing a low energy process compared to conventional thermal plasma conversion.

The problem of gas mixtures has been addressed in the past for CO₂ transport, focusing on the presence of small impurities and their influence in the thermophysical behavior of the CO₂ stream for pipeline design. In this work, the focus intends to expand the understanding of CCUS to assess the thermophysical behavior of gas mixtures involved in CO₂ decomposition through NTP in a fluidized plasma reactor (FPR). The electrons produced in the plasma acquire high energies (between 1 and 10 eV), and research shows that the use of argon (Ar) as a diluent gas proves beneficial in CO₂ conversion¹⁷ due to the lower minimum voltage required to form plasma with Ar. In effect, NTP decomposes CO₂ into carbon monoxide (CO) and oxygen (O₂) with the use of Ar as a diluent gas.¹⁷ Alternatively, other studies have shown the possibility to perform the process in a nitrogen (N₂) environment, resulting in an overall lower cost.¹⁸ Subsequently, the CO produced is usually coupled with a hydrogen source (H₂) to produce syngas (CH₄).¹⁹ In summary, mixtures of CO₂ with Ar, CO, O₂, N₂, H₂, and CH₄ can be found in NTP processes in a wide variety of compositions.

In this context, although extensive work has been devoted to the thermodynamic characterization of CO₂ mixtures, a review of the literature still identifies several gaps of knowledge, particularly when addressing either properties different than vapor–liquid equilibria (VLE) (i.e. second-order derivative properties) or when describing the behavior of multi-component mixtures. From a modeling perspective, the use of equations of state (EoSs), ranging from empirically specific-compound equations, classical and/or refined cubic equations, and statistical-mechanics-based equations, is common in the description of CO₂ mixtures in CCUS processes.

While it is not possible to review all works done in the area (the reader is referred to the extensive review of Li and co-workers²⁰ for further details), it is important to highlight some notable approaches. Among the semi-empirical EoSs developed for combustion gases, named EoS-CG, the equation of Span and Gernert,²¹ derived from the Span and Wagner EoS for pure CO₂,²² provides very high accuracy for the single and phase equilibria of mixtures with CO₂, CO, N₂, H₂, Ar, and H₂O at wide-ranging conditions. However, as any empirically-fitted equation, it suffers limitations when generalized to other components or conditions. Another pressure-explicit EoS studying these CCUS mixtures that has shown favorable results is the one published by Demetriades and Graham.²³ In terms of cubic equations, Peng–Robinson (PR) has been extensively used for these systems, completing experimental measurements, such as in the work of Vega-Maza and co-workers,²⁴ who provided new experimental data on the CO₂–H₂ and CO₂–N₂ mixtures, modeling them with PR. The work

of Carroll²⁵ also includes the PR equation, as well as the Soave–Redlich–Kwong (SRK) EoS, to study CO₂ mixtures with H₂S and CH₄. In addition, several authors provide a comparison between molecular-based statistical associating fluid theory (SAFT) equations and empirical or cubic thermodynamic models, as is the case for by Diamantonis et al.,²⁶ whose work provides a comparison between the SAFT,^{27,28} PC-SAFT,²⁹ and EoS-CG models for CO₂-rich multicomponent mixtures. Other SAFT models relevant in the context of CCUS (SAFT-VR Mie,³⁰ based on the Mie intermolecular potential) are also compared in another contribution.³¹

In this work, the polar version of soft-SAFT,³² an adequate thermodynamic model to treat the polarity in gases, is used to characterize the thermodynamic behavior of mixtures involved in CO₂ conversion processes with particular focus on the description of binary and ternary mixtures of common CCUS compounds relevant to NTP conversion: carbon dioxide (CO₂), argon (Ar), oxygen (O₂), carbon monoxide (CO), hydrogen (H₂), and nitrogen (N₂). Previous work shows that soft-SAFT has been successfully used to determine thermodynamic properties and phase equilibria calculations for other CO₂ mixtures involving a wide variety of compounds and solvents, such as polyether blends,³³ perfluoroalkanes,³⁴ monoethanolamine (MEA),³⁵ aqueous and non-aqueous amine hybrid solvents,^{36,37} ionic liquids,³⁸ and deep eutectic solvents,³⁹ as well as CH₄.⁴⁰

The present contribution is organized as follows: an initial overview of the theoretical background on soft-SAFT EoS is provided. Then, results for CCUS relevant compounds are presented and discussed: density behavior, derivative properties, second virial coefficients and VLE of CO₂ binary mixtures, thermodynamic VLE of other non-CO₂ binary mixtures, and, finally, the characterization of two ternary mixtures: CO₂, O₂, and Ar and CO₂, O₂, and N₂. Lastly, some concluding remarks are provided on the applicability of the soft-SAFT EoS in the context of CCUS.

2. THEORETICAL BACKGROUND AND METHODOLOGY

2.1. Soft-SAFT. Soft-SAFT^{41,42} is a well-known, mature variant of the original SAFT,^{27,28} a molecular-based EoS developed from statistical mechanics concepts that is capable of describing thermodynamic properties of complex fluids. It is obtained from Wertheim's first-order thermodynamic perturbation theory (TPT1) and accounts for different microscopic contributions such as high-range directional forces and chain formation.^{43–46}

As any SAFT model, soft-SAFT is expressed as the addition of residual Helmholtz energy of the system accounting for different microscopic contributions representing different molecular effects. For the particular case of non-associating ($A^{\text{assoc}} = 0$) polar molecules, such as those treated in this work, the equation reads:

$$A^{\text{res}} = A^{\text{ref}} + A^{\text{chain}} + A^{\text{polar}} \quad (1)$$

where A^{res} is the residual Helmholtz free energy ($A^{\text{total}} - A^{\text{ideal}}$), A^{ref} accounts for the free energy of reference of the component (understood as interaction between monomers), A^{chain} is the contribution due to chain formation, and A^{polar} accounts for component polarity. The reference term in soft-SAFT, A^{ref} , employs the Lennard–Jones (LJ) intermolecular potential for the repulsive and attractive interactions, computed following

the work of Johnson et al.⁴⁷ The expressions of the chain term, A^{chain} , resulting from Wertheim's theory (TPT1), is similar to other SAFT EoS versions. The contribution of multipolar interactions, A^{polar} , is obtained from the approach proposed by Gubbins and Twu,^{48,49} initially developed for spherical molecules and later extended to chains by Jog et al.⁵⁰ The perturbative theory used to develop the polar expression explicitly includes the polar moment (either the dipole, μ , or the quadrupole, Q), as well as another parameter representing the fraction of polar segments in a molecule (x_p). In this work, we are employing the interpolation equations for the integrals over pair and triplet correlation functions of a LJ fluid drawn from the work of Luckas et al.⁵¹ The reader is referred to the original soft-SAFT articles, including the subsequent addition of the polar contribution, for further details.^{32,41,42}

Based on the different terms previously shown, the description of the gases studied here within the soft-SAFT framework is given through a coarse-grained model defined by a homonuclear chain of m spherical segments of diameter σ_{ii} and dispersive energy between segments (ε_{ii}). In the case of quadrupolar fluids, the set of molecular parameters is completed with the effective quadrupolar moment (Q) and the fraction of segments containing the polar moment (x_p). This set of molecular parameters is sufficient to characterize a fluid and is usually obtained by fitting to experimental equilibrium data, most notably liquid density and vapor pressure data in a range of temperatures (from triple point, T_{tp} , to 0.95 times the critical temperature, T_c).

The extension of the theory to mixtures is straightforward for the case of the chain and polar terms as they are explicitly defined for mixtures. Concerning the reference LJ term, the segment diameter and different dispersive energies have to be averaged by generating a pseudo-compound with the thermodynamic properties of the mixture. This computational extension is achieved through the van der Waals one-fluid theory with the generalized Lorentz–Berthelot (LB) combining rules:

$$\sigma_{ij} = \eta_{ij} \left(\frac{\sigma_{ii} + \sigma_{jj}}{2} \right) \quad (2)$$

$$\varepsilon_{ij} = \xi_{ij} \sqrt{\varepsilon_{ii} \varepsilon_{jj}} \quad (3)$$

where σ_{ij} and ε_{ij} are the crossed segment diameter and dispersive energy for the mixture, and η_{ij} and ξ_{ij} are the adjustable size and energy binary parameters to account for deviations from the LB rule. These parameters consider the differences in size and/or energy between the monomers from the compounds in the mixture, and they are commonly fitted to isothermal or isobar equilibrium data.

2.2. Molecular Models. In this work, six gases have been modeled: CO₂, CO, N₂, O₂, H₂, and Ar as indicated in Table 1.

Table 1. Chemical Compounds and Models

chemical name	chemical formula	CAS number	model
carbon dioxide	CO ₂	124-38-9	soft-SAFT
carbon monoxide	CO	630-08-0	soft-SAFT
nitrogen	N ₂	7727-37-9	soft-SAFT
oxygen	O ₂	7782-44-7	soft-SAFT
hydrogen	H ₂	1333-74-0	soft-SAFT
argon	Ar	7440-37-1	soft-SAFT

All are described as homonuclear chains of m spherical segments of diameter σ and monomer–monomer dispersive energy ε/k_B . The quadrupole effect in CO₂, CO, and N₂ has been explicitly considered. In this case, the polarity fraction is fixed a priori, and the quadrupole moments are obtained from experimental values.³⁶ As in previous studies,^{52,53} dipole moment of CO is not taken into account in the molecular characterization due to its very low value ($\mu = 0.122$ Debye).⁵⁴ In the case of O₂ and H₂, their very low polarity is considered negligible for their characterization and no quadrupole moment is explicitly expressed. Concerning Ar, an apolar spherical molecule, an additional constraint is imposed by fixing the chain length (m) to 1. The final set of optimized molecular parameters for all compounds is taken from previous contributions, with the exception of Ar, which has been refitted using saturated liquid density and vapor pressure data to improve its phase equilibrium description. Table 2 summarizes the molecular parameters of all the compounds involved in this work.

The phase equilibrium of the thermodynamic behavior of Ar is presented in Figure S1 of the Supporting Information. Excellent agreement is found between the experimental data and the soft-SAFT description with an average absolute deviation (AAD%) of 0.717 and 1.076% for the saturated liquid density and vapor pressure, respectively.

3. RESULTS AND DISCUSSION

3.1. Binary Mixtures Characterization. Once the parameters for the pure gases are well established, binary mixtures are characterized. The CO₂ binary mixtures modeled in the present work are mixtures of CO₂ with Ar, O₂, CO, N₂, and H₂. Additional binary mixtures without CO₂ include O₂ and Ar, CO and N₂, and CO and Ar. A summary of all calculations carried out for all mixtures, including all properties and operating conditions (temperature, pressure, and composition) are summarized in Table S1 in the Supporting Information.

First, the single-phase density of several CO₂ mixtures is predicted at different temperatures, ranging from the subcritical till the supercritical region, as shown in Figures 1 and S3a (for CO₂–H₂) in the Supporting Information.

The comparison with experimental data reveals very good agreement for all the mixtures studied with some minor deterioration at the highest pressures and lowest temperatures for the CO₂–Ar system. The information provided by the equation is useful given the experimental gaps noticed for several mixtures. In particular, $p\rho T$ experimental data are scarce for the CO₂–CO mixture,⁵⁸ of particular interest in this work, as it only covers a small phase region.

Apart from the effect of the pressure on the density, it is desirable to analyze additional properties to check the validity of the parametrization. The use of derivative properties is always a good test as they are more sensitive to deviations than primary properties. In Figure 2, the speed of sound at different isotherms is predicted as a function of pressure for the CO₂–Ar mixture at two given compositions. The speed of sound can be estimated from the combination of other derivative properties, such as the isochoric and isobaric heat capacities, according to:

$$\omega = \sqrt{\frac{C_p}{C_v} \left(\frac{1}{\rho \times k_T} \right)} \quad (4)$$

Table 2. Optimized Soft-SAFT Molecular Parameters for All Pure Compounds

compound	m	σ (Å)	ε/k_B (K)	$Q \times 10^{40}$ ($C \times m^2$)	x_p	reference
CO ₂	1.571	3.166	166.5	14.68	0.333	40
CO	1.156	3.500	85.5	9.473	0.500	40
N ₂	1.304	3.291	85.2	4.53	0.500	40
O ₂	1.168	3.198	111.5			55
H ₂	0.487	4.244	33.9			55
Ar	1.000	3.401	117.1			this work

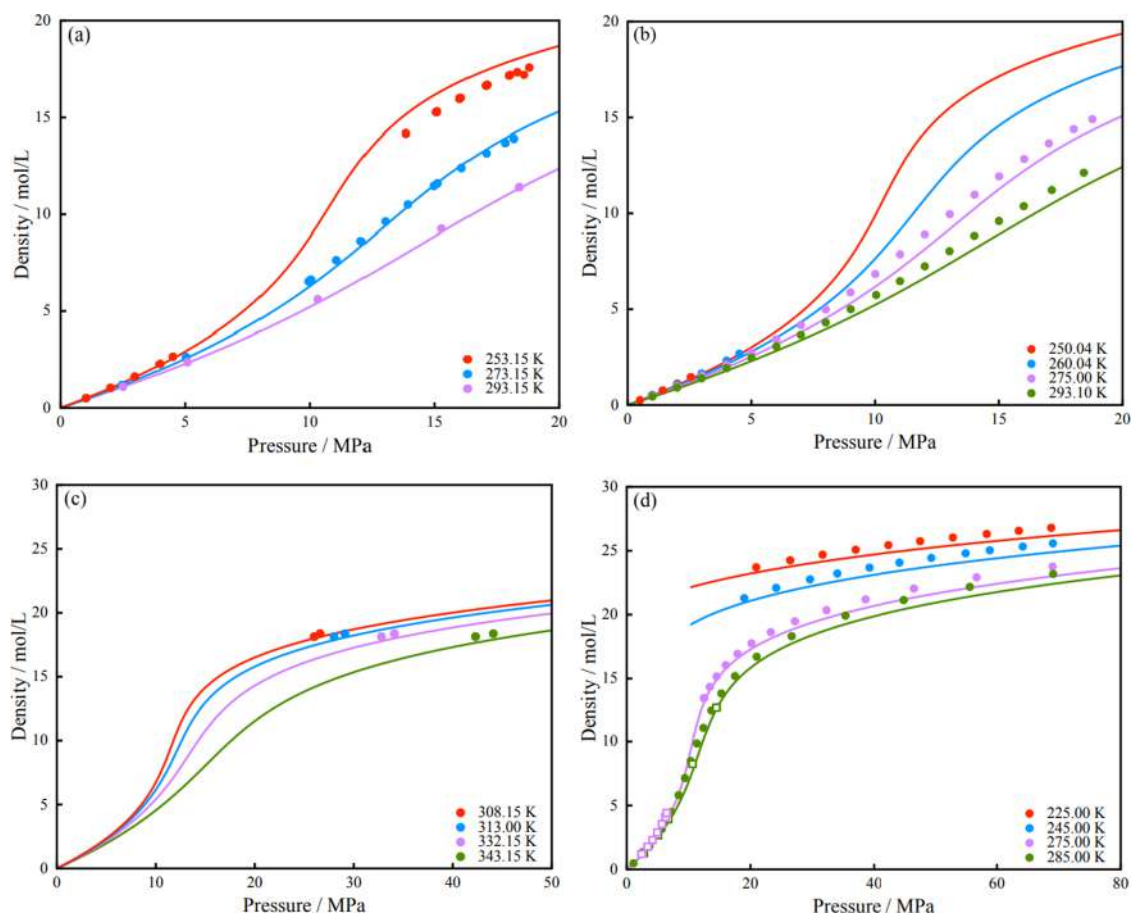


Figure 1. Isothermal density–pressure diagrams of CO₂ mixtures with (a) 50% Ar with experimental data (circles⁵⁶), (b) 50% O₂ with experimental data (circles⁵⁷), (c) 13.8% CO with experimental data (circles⁵⁸), and (d) 28.895% N₂ with experimental data (squares⁵⁹ and circles⁶⁰). In all cases, soft-SAFT calculations are represented by full lines.

where c_p and c_v are the isobaric and isochoric heat capacities, respectively, k_T is the reduced bulk modulus, and ω is the speed of sound. In these calculations, the ideal contribution to the heat capacities is added from the available correlations given at NIST.⁶¹ Good agreement is found in all cases, with a very slight over prediction at the lowest temperatures. Still, the qualitative prediction is excellent, given the sensitivity of this property.

Given the fact that the speed of sound has been compared at isotherms which are at a temperature higher than the working conditions of interest, other properties have been evaluated so as to further validate the approach. In Figure 3a, the isothermal compressibility as a function of pressure of the CO₂–Ar mixture has been predicted in a wide range of temperatures. In this system, the predictions underestimate the isothermal compressibility values although good qualitative behavior is achieved in all cases.

In addition, the prediction of the second virial coefficient for the same mixture (CO₂–Ar), which is a reliable assessment to check the validity of the model, has been included. Notice that the green line represents pure argon, showing an excellent agreement with the data. The second virial coefficient of the mixtures is also in good agreement with the experimental data with some minor deviations that are constant in the whole temperature range of study (250–350 K).

Next step concerns the evaluation of the binary VLE of these systems. The study of the VLE of CO₂–gas mixtures has already been described by other SAFT-type equations as mentioned in the literature.^{26,31,65} The nature of these mixtures does not allow quantitative agreement without further adjustment of binary parameters, a common characteristic in all previous bibliography efforts. When using soft-SAFT, a similar performance was found (prediction using combining rules without any binary parameter deviates from the experimental data as shown in Figure S2 of the Supporting Information for the mixture

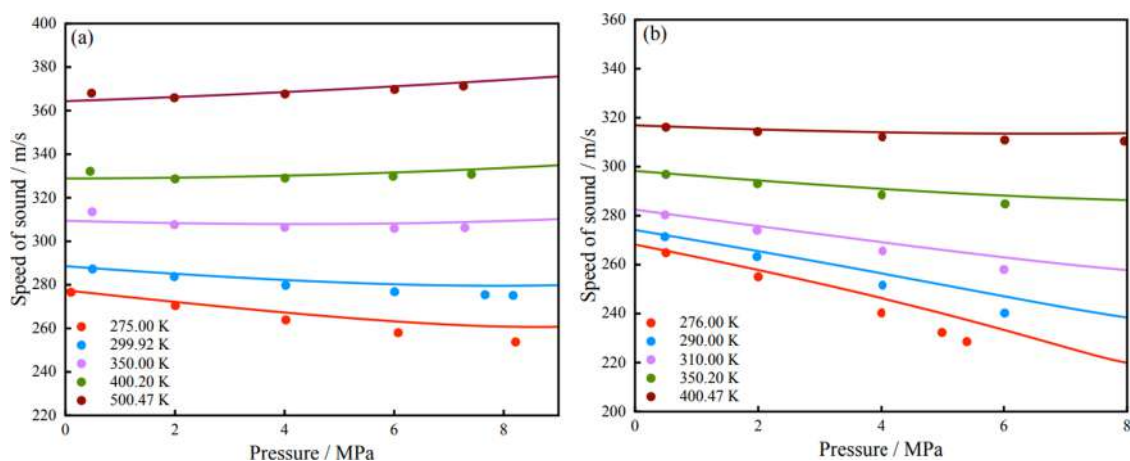


Figure 2. Isothermal speed of sound–pressure diagrams for CO₂ and Ar mixture with CO₂ compositions of (a) $x_{\text{CO}_2} = 0.50104$ and (b) $x_{\text{CO}_2} = 0.74981$. In all cases, symbols⁵⁶ are experimental data and full lines represent soft-SAFT predictions.

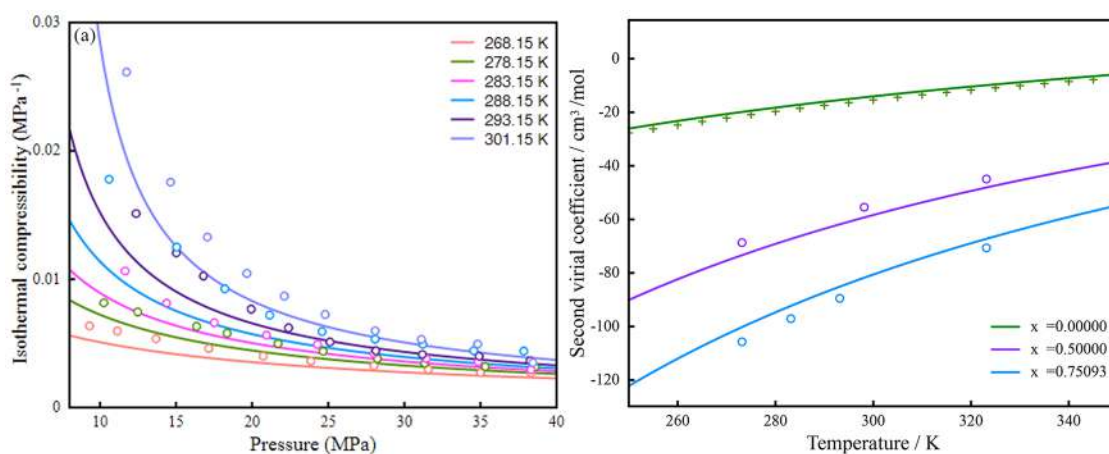


Figure 3. Predicted properties for the CO₂ + Ar mixture. (a) Isothermal Compressibility of CO₂ with Ar at $x_{\text{CO}_2} = 0.0446$ and different temperatures. Symbols represent experimental data.⁶⁴ (b) Second virial coefficients of pure Ar and a mixture of CO₂ (1) with Ar (2) as a function of temperature at different compositions. Crosses,⁶² and circles,⁶³ represent experimental data. In both Figures, lines are the soft-SAFT predictions.

CO₂–Ar). This is a consequence of the incondensable nature of the gases involved and the arbitrary choice of the combining rules. Consequently, and following the same procedure done by other authors,^{26,31,65} a correcting size and energy binary parameters (η and ξ , respectively) are introduced to fit the results at all isotherms using available experimental data. In this work, the VLE representation of all the mixtures is achieved with a constant binary energy parameter (ξ) close to unity, while a binary temperature-dependent size parameter η has been fitted and correlated with a second-order polynomial form for the CO₂–gas mixtures. The only exception is the CO₂–H₂ mixture, where a good description for all mixtures is achieved without using any size parameter.

Although it seems more common to establish an energy temperature-dependent binary parameter, we have noticed that a size binary parameter can better capture the VLE behavior of gas mixtures due to entropic effects as a consequence of the relatively low energy interaction with incondensable gases such as Ar, N₂, CO, or H₂. In order to demonstrate this sensitivity of the η binary parameter on the VLE description, a complete characterization study of the impact of both binary parameters, η and ξ , has been carried out using the CO₂–Ar system at the benchmark system. Figure 4 shows a parametric analysis of the carbon dioxide plus argon system at four temperatures. The

average deviation (%) contrasted with the experimental data is displayed as contour plots.

In Figure 4, each projection shows summit and valley areas, particularly at low temperatures, caused by the emergence of equilibrium features, such as liquid–liquid equilibrium (see Figure 4a). In all cases, the global minimum is highlighted with (A) for each temperature and lies in a coordinate where both interaction coefficients are different than one. Points (X) are the interaction parameters selected in the study using the presented correlation. In two cases, Figure 4b,d, point (A) matches with point (X). Additionally, at all temperatures, the ADD of points (A) and (X) are <1.60% and <3.00%, respectively. Other valley zones appear as in Figure 4b, highlighted in point (B). These local optima are discarded due to the underprediction of the mixture critical point.

The most relevant conclusion of the study is the fact that while the ξ dispersive interaction coefficient has a stable value of around 1.01 in all cases, the η optimum is more affected by the temperature change. More importantly, it can be seen that the blue areas (indicating the lowest deviations) cannot be achieved using an η parameter equal to unity. Indeed, the η parameter becomes significantly higher than 1 at the highest temperatures. For this reason, a temperature functionality in the η parameter has been used.

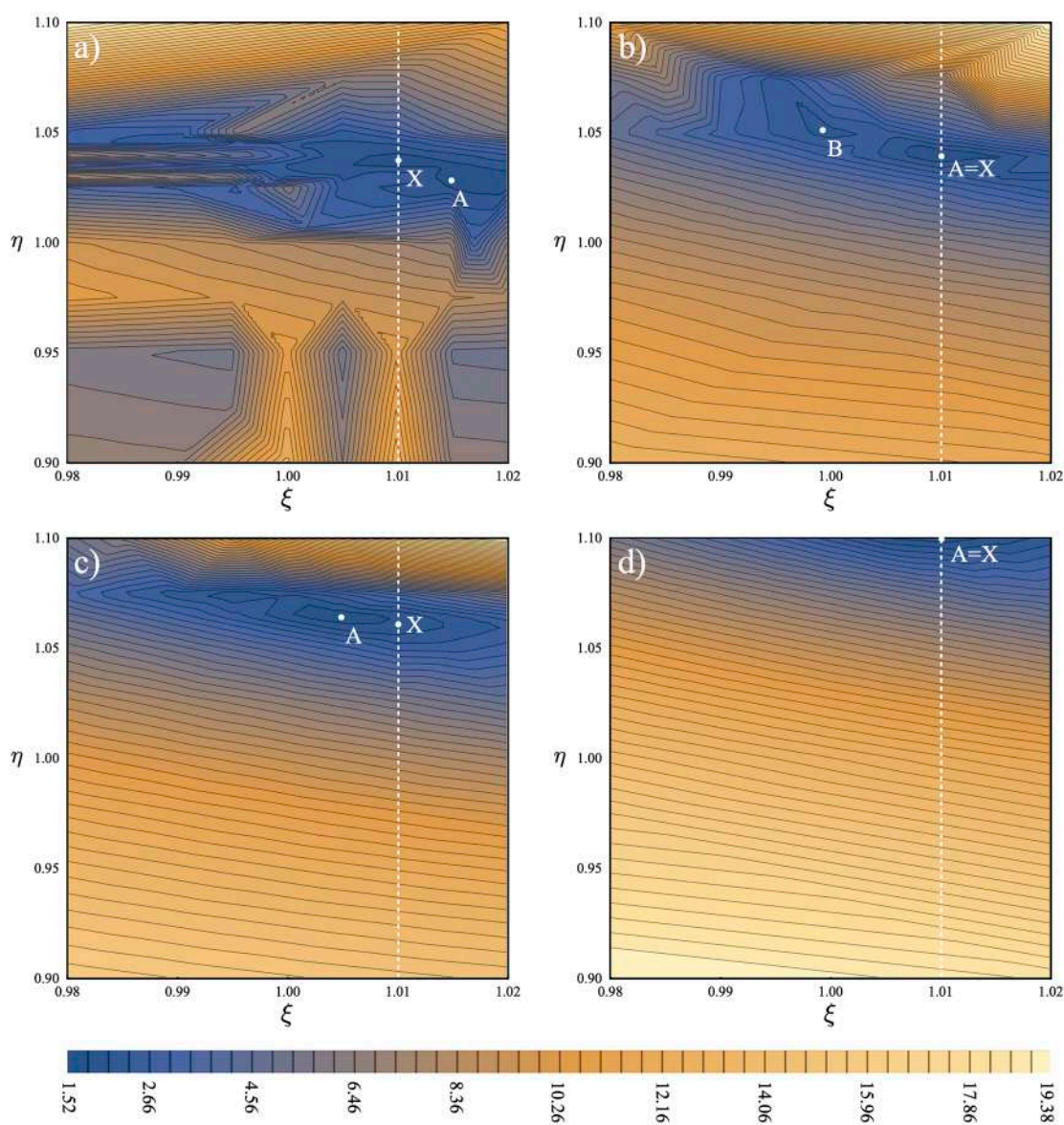


Figure 4. Parametric analysis of the influence of binary parameters ξ and η in the description of the vapor–liquid equilibria of CO₂ with Ar at (a) 223.15 K, (b) 233.32 K, (c) 253.28 K, and (d) 273.26 K. Letter A indicates the minimum absolute average deviation (%AAD), letter B shows a local minimum, and letter X indicates the binary parameter value selected.

A summary of the required coefficients and values for the η and ξ binary parameters, respectively, is given in Table 3. The study of η correlation to temperature allows the extrapolation to other temperatures when lacking experimental data. Still, it

Table 3. Summary of Size (η) and Energy (ξ) Binary Parameters^a

mixture	η			ξ
	$a \times 10^5$	$b \times 10^3$	c	
CO ₂ –Ar	2.06206	−8.86548	1.98277	1.010
CO ₂ –O ₂	1.12766	−4.45338	1.45999	1.010
CO ₂ –CO	1.00000	−4.16299	1.47101	1.130
CO ₂ –N ₂	1.58793	−6.96290	1.79146	1.100
CO ₂ –H ₂	0	0	1.00000	1.350

^a η follows a second-order degree polynomial with temperature according to the expression: $aT^2 + bT + c$.

is significant to note that limitations on the usage of these correlations can appear when approaching critical boundaries of the mixtures.

The final results of the VLE modeling are displayed in Figures 5 and S3b, (for the CO₂–H₂ mixture) in Supporting Information, showing several isotherms for all mixtures (ranging from 218.15 to 293.15 K). As it can be noticed, very good agreement is achieved between the soft-SAFT model (lines) and the available experimental data (symbols). Although some minor discrepancies are observed in the vicinity of the critical point, particularly over ambient temperature ($T > 273.15$ K), the overall representation provides a reliable characterization of the VLE behavior (Figure 5).

The comparative analysis among the different CO₂–gas mixtures allows better understanding on how they thermodynamically behave. For this purpose, the comparison of five mixtures (including CO₂–H₂) at the same isotherm is

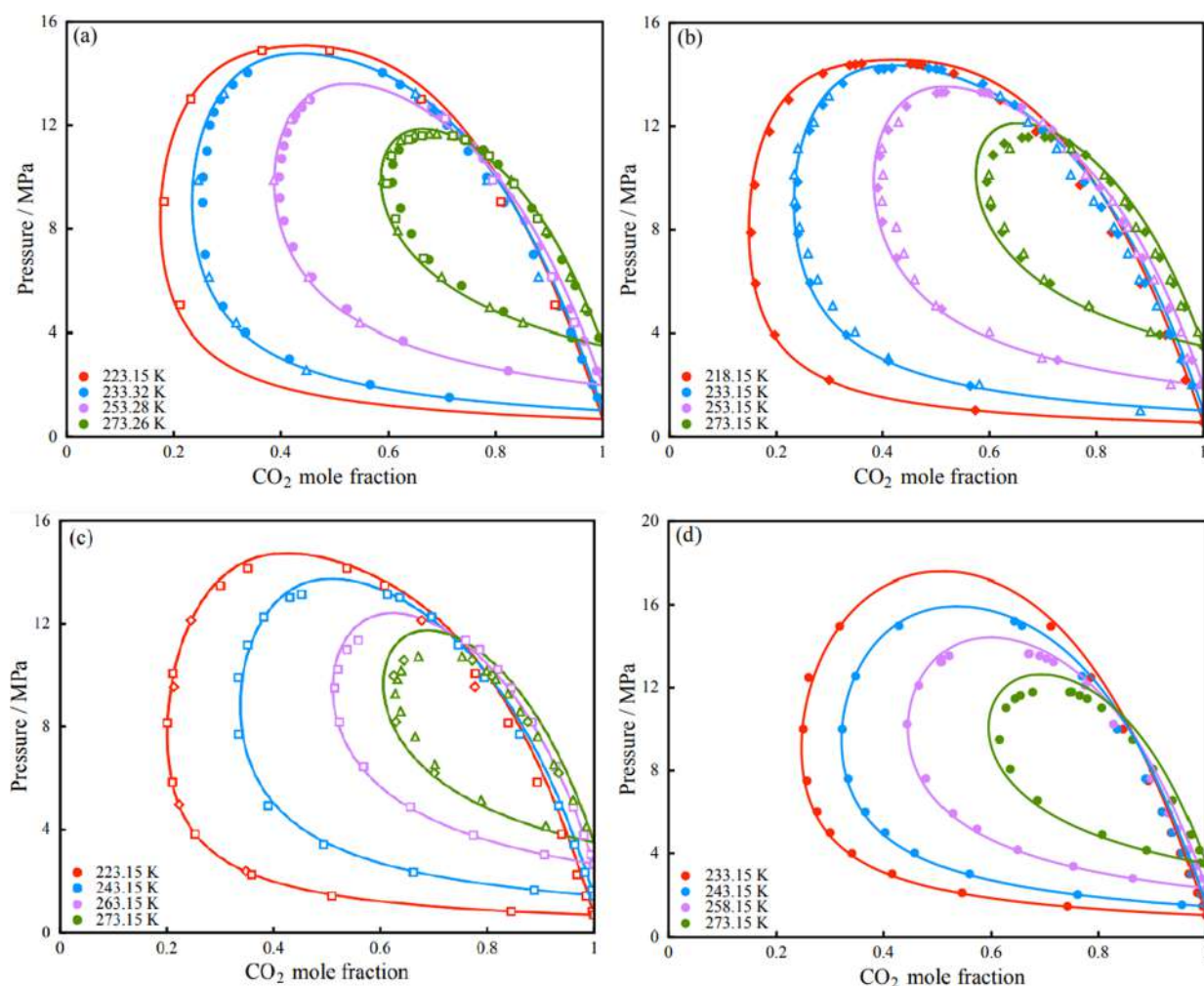


Figure 5. Isothermal vapor–liquid equilibrium of CO₂ at four temperatures with (a) Ar with experimental data (filled circles,⁶⁶ squares,⁶⁷ and triangles⁶⁸), (b) O₂ with experimental data (triangles⁶⁹ and filled diamonds⁷⁰), (c) CO with experimental data (squares,⁷¹ diamonds,⁶⁸ and triangles⁷²), and (d) N₂ with experimental data (circles²⁴). In all cases, soft-SAFT calculations are represented by full lines.

presented in Figure 6. While several mixtures behave similarly at higher concentrations of CO₂ in the vapor phase, remarkable variations are observed in the liquid phase and, consequently, in the phase envelope. While the effect of Ar, O₂, and CO is very similar, the addition of nitrogen notably increases the vapor pressure of the system. Finally, the impact of H₂ is dramatic due to the presence of small hydrogen molecules in the system. It is interesting to note that these differences can be observed in the values of the soft-SAFT molecular parameters included in Table 2.

In this work, other prevalent binary mixtures not containing CO₂ have also been characterized with the soft-SAFT EoS, considering they can be found in NTP conversion processes. A major limitation in these non-CO₂ mixtures is the limited availability of VLE data, particularly in the case of binary mixtures with carbon monoxide as it occurs for the CO–O₂ and CO–H₂ mixtures, where no experimental data have been found. This data scarcity is possibly due to the difficulty in working at very low temperatures as a consequence of the low critical points of these gases.

The description of the O₂–Ar and CO–N₂ mixture in the range 68–100 K is given in Figure 7. It is notable to remark that the absence of CO₂ increases the ideality of the system and, consequently, excellent agreement can be found using a

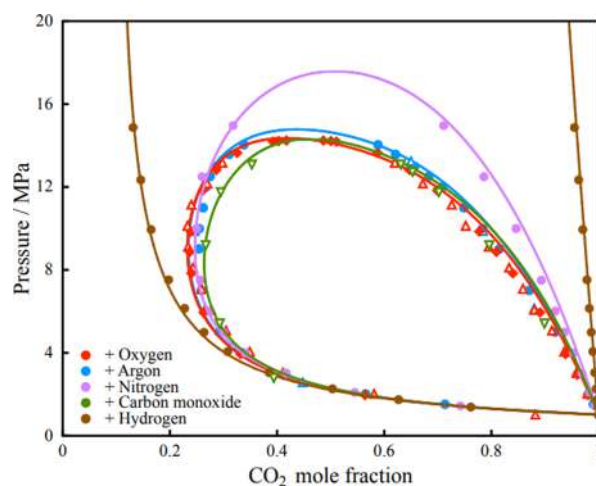


Figure 6. Isothermal vapor–liquid equilibrium of CO₂ mixtures at 233 K with Ar (experimental data: circles⁷³), O₂ (experimental data: triangles⁶⁹ and diamonds⁷⁰), CO (experimental data: inverted triangles⁷¹), N₂ (experimental data: circles²⁴), and H₂ (experimental data: circles⁷⁴). Soft-SAFT calculations are represented by full lines.

single ξ binary parameter (while η is set to unity), which also turns out to be very close to 1.00 in most cases (O₂ and Ar: ξ =

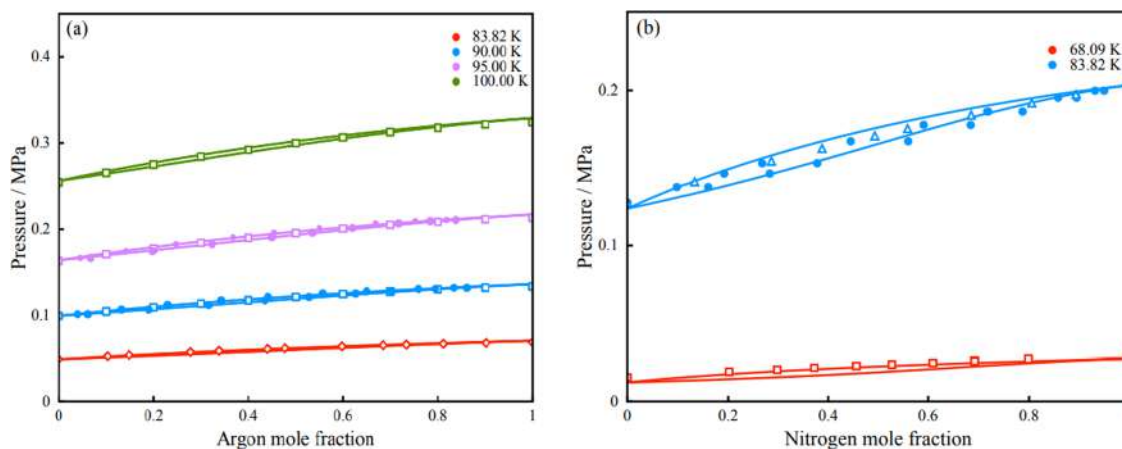


Figure 7. Isothermal vapor–liquid equilibrium of (a) O_2 and Ar mixture with experimental data (squares,⁷⁵ filled circles,⁷⁶ and diamonds⁷⁷), and (b) CO and N_2 with experimental data (triangles,⁷⁷ filled circles,⁷⁸ and squares⁷⁹). In all cases, soft-SAFT calculations are represented by full lines.

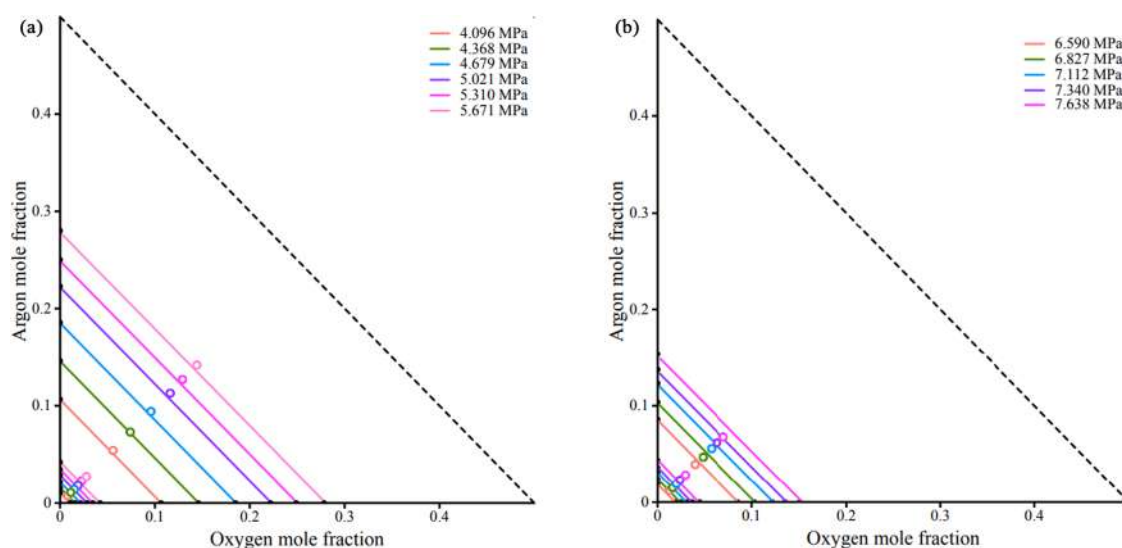


Figure 8. Ternary diagram of CO_2 , O_2 , and Ar mixture at (a) $T = 273.24$ K and (b) $T = 293.21$ K. In all cases, symbols are experimental data (circles⁷³) and soft-SAFT predictions are represented by full lines.

0.99, CO and N_2 ; $\xi = 1.01$). Additionally, the mixture of CO –Ar ($\xi = 1.06$), also with η set to unity, can be found in Figure S4 in the Supporting Information.

3.2. Ternary Mixture. One of the main goals of this characterization is to be capable to reproduce multicomponent mixtures. Curiously, although the interest for these mixtures is evident and has practical applications, the amount of work devoted to the description and modeling of ternary systems decreases exponentially compared to binary systems. The previous characterization should enable soft-SAFT to estimate ternary data. Considering that CO_2 , O_2 , and Ar or N_2 (depending on the type of diluent gas used) are the compounds with the highest concentrations in NTP CO_2 conversion processes, ternary mixtures of CO_2 , O_2 , and one of these inert gases are particularly important for the current study.

First, the ternary equilibrium of this system (CO_2 , O_2 , and Ar) is predicted at three different temperatures and several isobars, and compared to experimental data sets,⁷³ as depicted in Figures 8 and S5 (Supporting Information). In all cases, binary parameters are transferred for each binary pair without further adjustment at 273.24 K, while the parameter

correlations are used, when necessary, to predict the behavior at 293.21 K. In both cases, soft-SAFT described the vapor–liquid phase boundaries of the system in good agreement with the experimental data with slight deviations in the vapor phase that increase with the pressure. Still, the results are satisfactory, showing the predictive capacity of the equation. As an additional remark, it is interesting to note that the predictions of these systems without any binary parameter, done for comparison in Figure S6 (Supporting Information), deteriorate at the highest temperature in an equivalent manner as it occurs with the binary systems. These deviations are expected to become more significant at higher pressure if no binary interactions are considered.

In order to extend the validation to other ternary systems, we have also studied the ternary mixture involving CO_2 – O_2 – N_2 as shown in Figure 9. In this case, the predictions are in excellent agreement with the experimental data for the liquid phase at the two studied temperatures, reaching also a good description of the vapor phase. The quasi-linear behavior of the isobars, already observed in Figure 8, is repeated here.

Finally, it is important to remark that the prediction without binary parameters (see Figure S7 in the Supporting

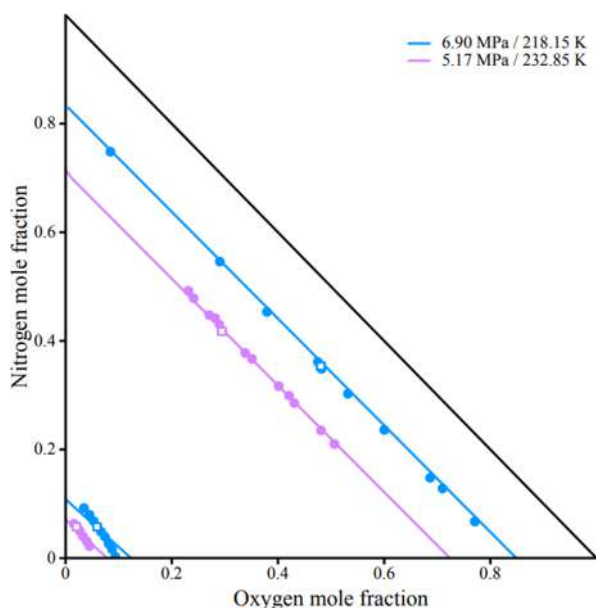


Figure 9. Ternary diagram of CO₂, O₂, and N₂ mixture at $T = 218.15$ K (blue) and 232.85 K (violet). In all cases, symbols are experimental data (circles⁸⁰ and squares⁸¹) and soft-SAFT predictions are represented by full lines.

Information) reveals severe deviations in the liquid phase. This is the result of an incorrect prediction of the phase behavior of the CO₂–N₂ system without binary parameters, where a Type III behavior is predicted and, consequently, the liquid phase is found at higher pressure. This reinforces the idea that it is necessary to account for the binary interactions so as to ensure a proper description of multicomponent systems.

4. CONCLUSIONS

In this work, the modeling of single-phase and phase equilibria of gas mixtures has been described through soft-SAFT EoS in the context of CO₂ conversion. In particular, a thermodynamic study of the compounds and common impurities involved in NTP CO₂ conversion processes (CO₂, CO, Ar, O₂, N₂, and H₂) has been carried out.

The molecular parameters of all relevant pure substances have been stated from previous literature, except in the case of Ar. These molecular models consider the physical meaning of the molecules, including their polar contributions. In the case of binary mixtures, fitting to experimental data has been achieved through adjustable size and energy binary parameters. A complete parametric analysis has demonstrated the temperature dependency of the size binary parameter exhibited by these gas mixtures. The thermodynamic characterization of each mixture has been completed for a range of conditions with available VLE, including the prediction of $p\rho T$, derivative properties, and second virial coefficients. All mixtures have shown good fitting to experimental data, albeit some limitations appearing when nearing critical regions, as expected. For CO₂-rich mixtures, the correlation found between temperature and binary parameters allows the study of all mixtures in a predictive manner in the case where no experimental data are available at a specific isotherm.

Finally, the study of two ternary mixtures (CO₂ and O₂ with Ar or N₂) has been provided in a fully predictive manner by transferring the binary parameters obtained for the binary systems or using the correlations proposed for these

parameters. Results have shown an accurate representation of the multicomponent mixtures' behavior, opening the door for further multicomponent calculations, where data are scarce. These results demonstrate the usefulness of SAFT-type equations, where a balance between complexity and accuracy is achieved.

■ ASSOCIATED CONTENT

SI Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.jced.3c00131>.

Table S1: Summary of all thermodynamic properties, temperature, pressure, and composition studied in this work; Figure S1: Thermodynamic characterization of argon (Ar) with soft-SAFT; Figure S2: Predicted isothermal vapor–liquid equilibrium of CO₂ with Ar at four temperatures; Figure S3: Isothermal behavior for CO₂ and H₂ mixture at selected temperatures modeled with soft-SAFT (a) density–pressure diagram (b) vapor–liquid equilibrium; Figure S4: Vapor–Liquid equilibrium of CO and Ar mixture modeled with soft-SAFT; Figure S5: Predicted ternary diagram of CO₂, O₂, and Ar mixture at (a) $T = 253.27$ K; Figure S6: Predicted ternary diagram of CO₂, O₂, and Ar mixture at (a) $T = 273.24$ K and (b) $T = 293.21$ K at different isobars with and without transferring binary parameters; Figure S7: Predicted ternary diagram of CO₂, O₂, and N₂ mixture at $T = 218.15$ K (blue) and 232.85 K (violet) (PDF)

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Notes

The authors declare no competing financial interest.

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